



R&DE (Engineers), DRDO

Unsymmetrical Bending of Beams

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Introduction

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- Beam – structural member – takes transverse loads
- Cross-sectional dimensions much smaller than length
- Beam width same range of thickness/depth
- Thin & thick beams
- If $l \geq 15 t$ – thin beam
- Thin beam – Euler – Bernoulli's beam
- Thick beam – Timoshenko beam



Introduction

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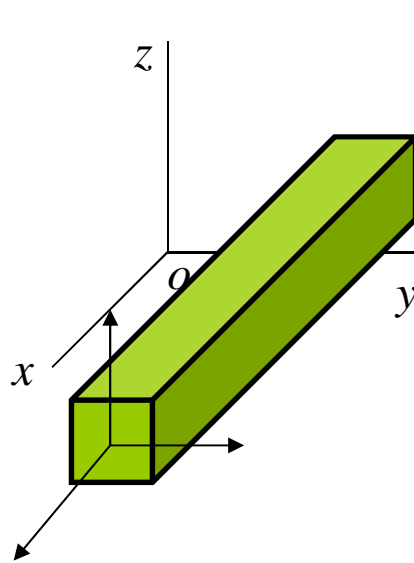
- In thin beams – deformation due to shear negligible
- Thick beams – shear deformation considered
- Beam – one dimensional structural member – length is very high than lateral dimensions
- Various parameters => function of single independent variable



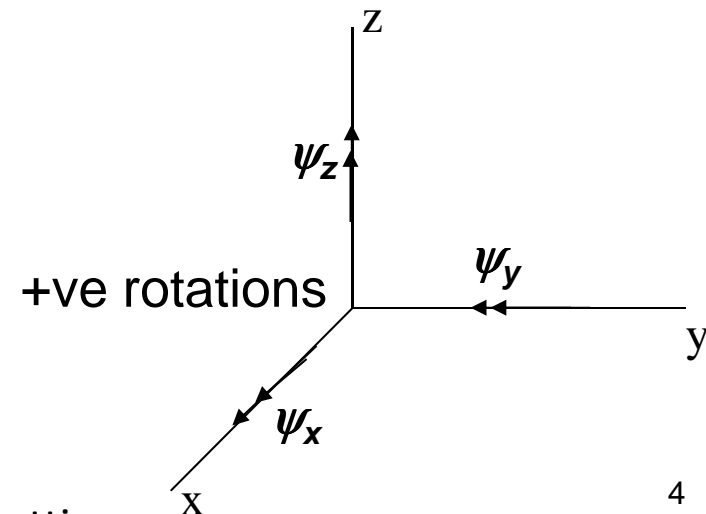
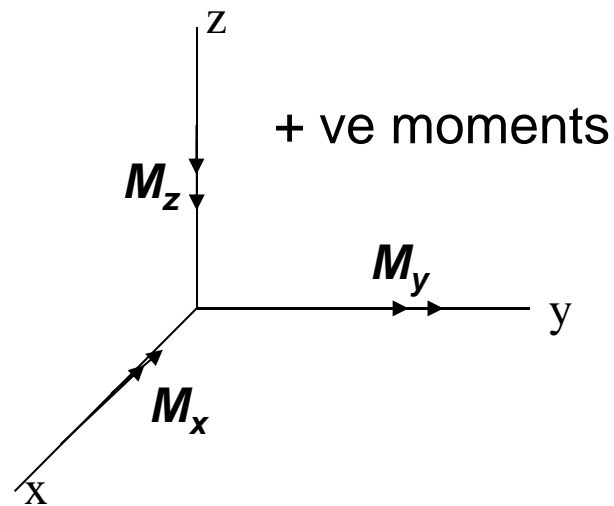
Sign convention

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- Following sign convention is followed



Orientation of beam



Beam bending in xy and xz planes

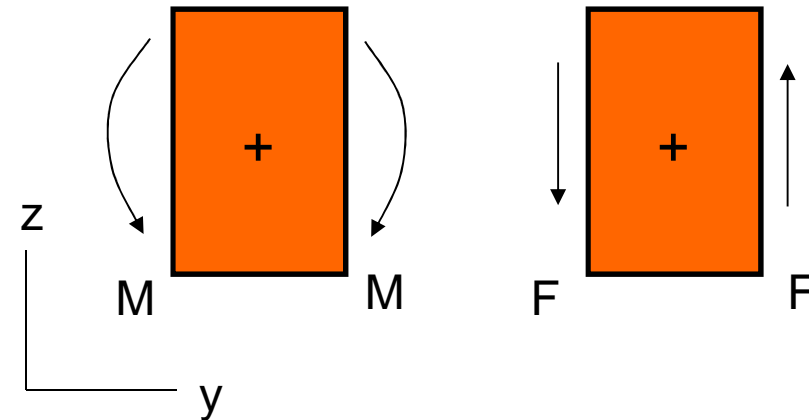
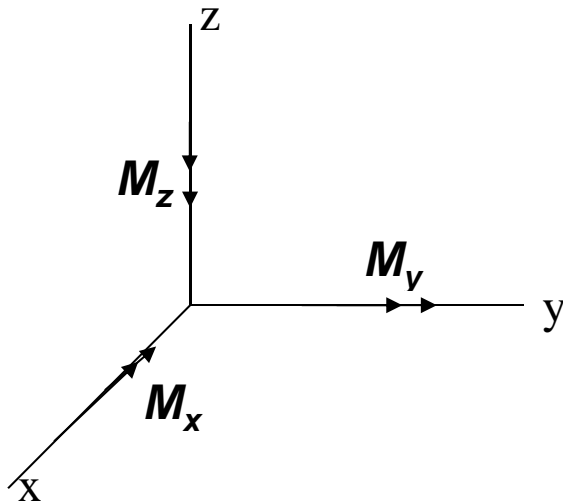
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Sign convention

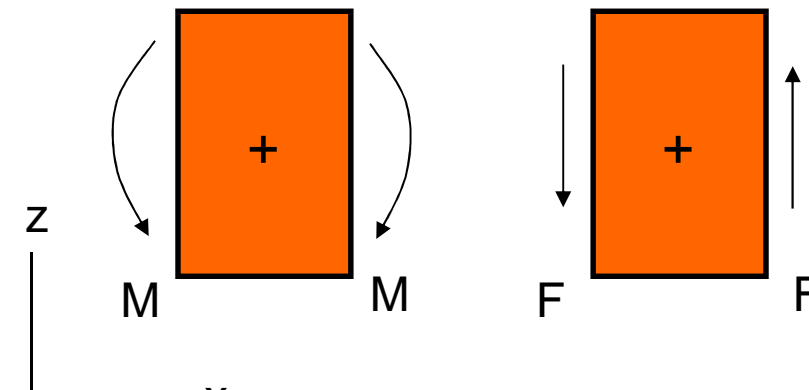
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■ Shear force and bending moments



Positive shear force and bending moments

Convex upward +ve direction



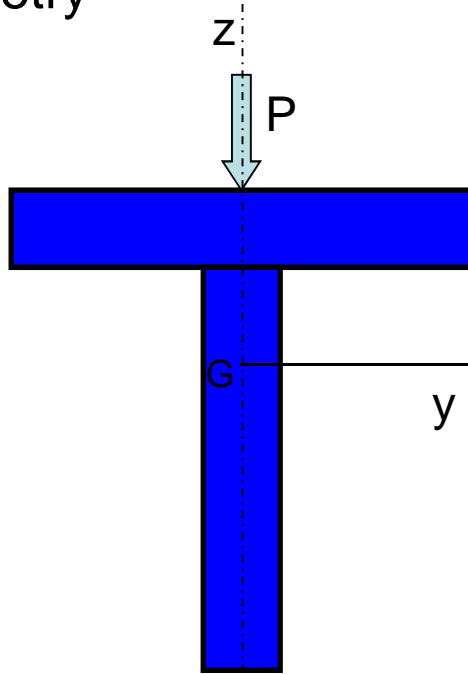
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Unsymmetrical bending

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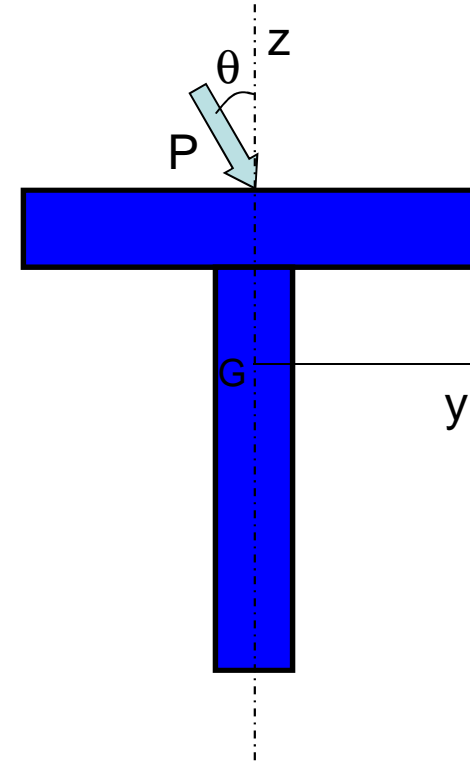
Load applied in the plane of symmetry



Symmetrical cross-sectional – load applied in the plane of symmetry - xz

Bending takes place in xz plane

Load applied at some orientation



Bending takes place in both planes – xz and xy

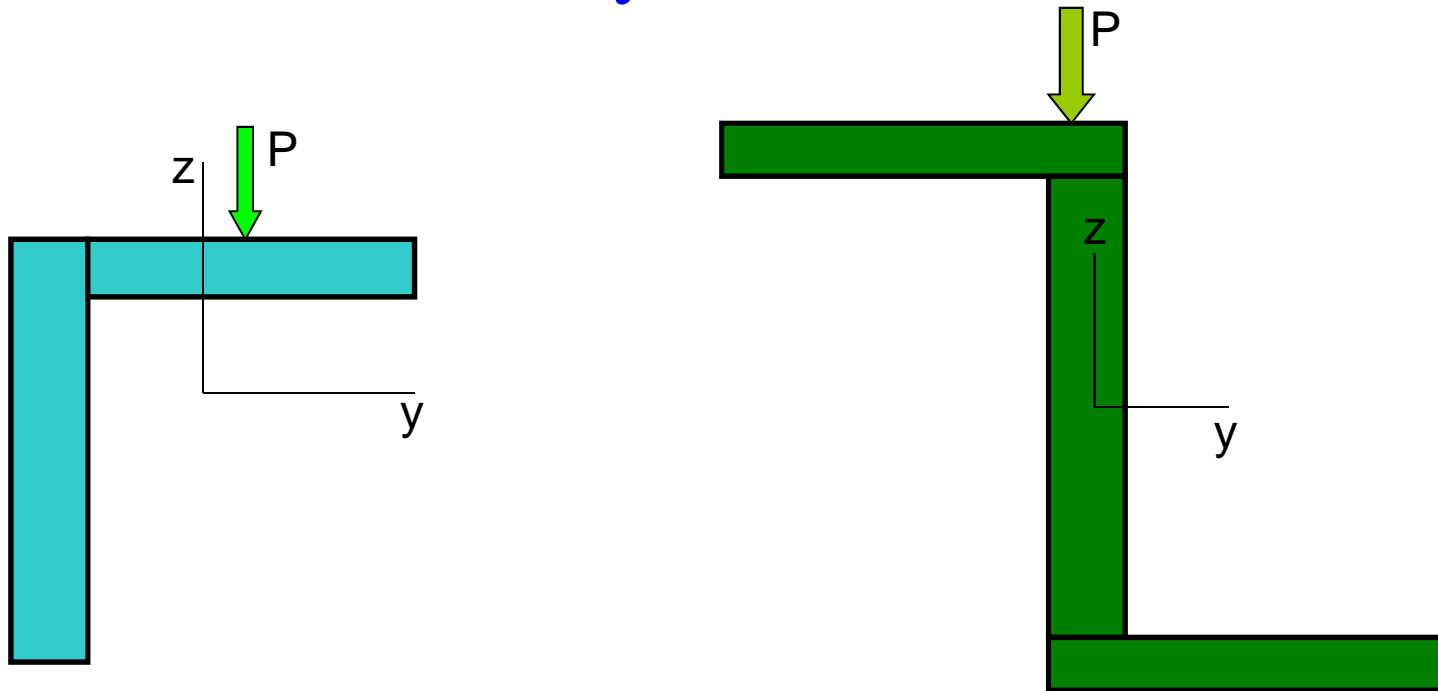
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Unsymmetrical bending

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- Cross-section is unsymmetrical



Bending takes place in both planes

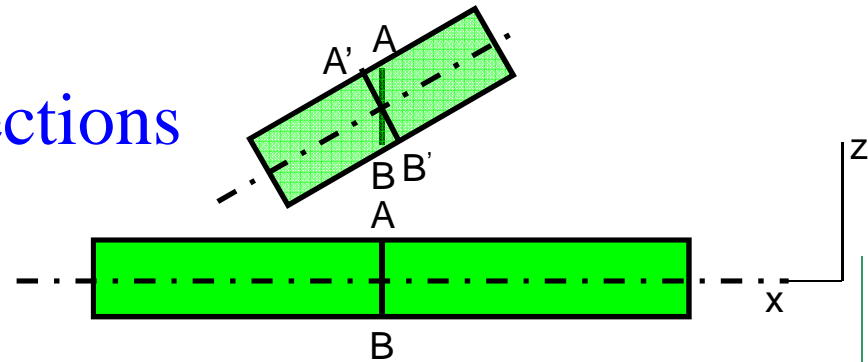
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Euler – Bernoulli's beam theory

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- Basic assumptions
 - Length is much higher than lateral dimensions
– $l \geq 15 t$
 - Plane cross section remains plane before and after bending
 - Stresses in lateral directions negligible
 - Thin beam
strain variation is linear across cross-section
 - Hookean material – linear elastic



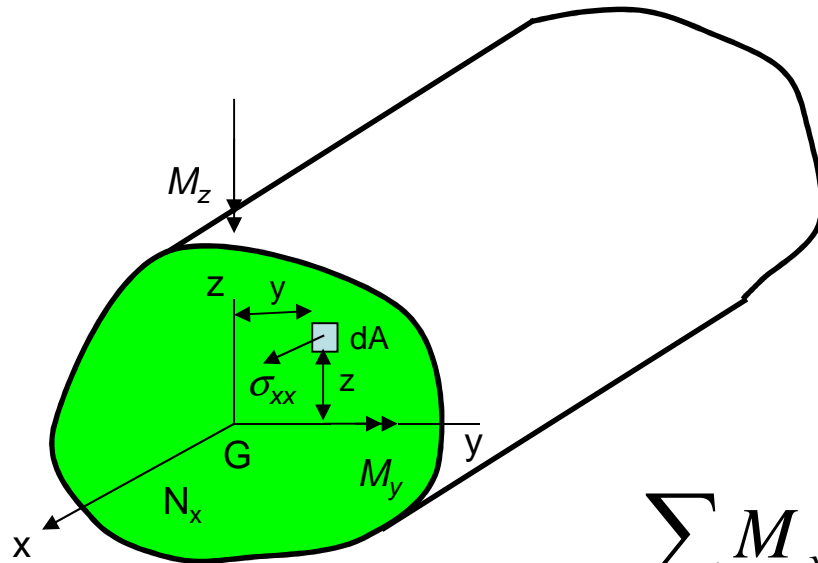
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Bending of arbitrary cross section beam

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- An arbitrary cross-section beam oriented along x -direction



xyz csys coincides with centroid

$$\epsilon_{xx} = ay + bz + c$$

$$\sigma_{xx} = E\epsilon_{xx} = E(ay + bz + c)$$

Equilibrium of the section requires

$$\sum F_x = N_x$$

$$\sum M_y, \sum M_z \Rightarrow \text{External moments}$$

$$dN_x = \sigma_{xx} dA = E(ay + bz + c)dA$$

Integrate this for total force in x - direction

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Bending of arbitrary cross section beam

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- This equals to external force in 'x' direction = N_x

$$N_x = \int_A \sigma_{xx} dA = \int_A E(ay + bz + c) dA$$

$$N_x = Ea \int_A y dA + Eb \int_A z dA + Ec \int_A dA$$

$$N_x = EcA \Rightarrow c = \frac{N_x}{EA}$$

Origin of xyz co-ordinate system is selected at centroid

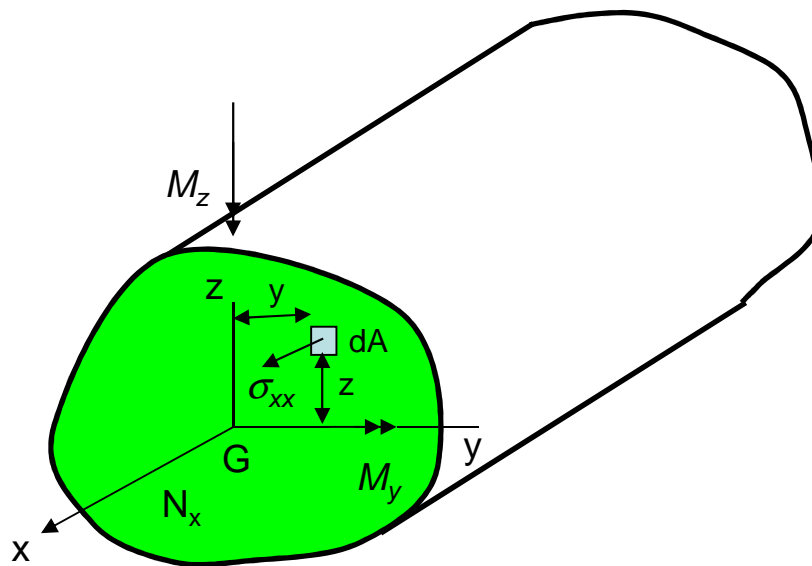
- Coefficients of 'Ea' and 'Eb' constants vanish



Bending of arbitrary cross section beam

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■ Moment wrt 'y' axis



$$dM_y = (\sigma_{xx} dA) z$$

$$dM_y = Ez \varepsilon_{xx} dA$$

Integrating over cross-sectional area 'A'

$$M_y = \int_A Ez (ay + bz + c) dA$$

$$\Rightarrow M_y = Ea \int_A yz dA +$$

$$Eb \int_A z^2 dA + Ec \int_A z dA$$



Bending of arbitrary cross section beam

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■ Moment about 'y' axis

$$M_y = Ea \int_A yz \, dA + Eb \int_A z^2 \, dA + Ec \int_A z \, dA$$
$$\Rightarrow M_y = Ea I_{yz} + Eb I_{yy} \quad - (1)$$

■ Moment about 'z' axis

$$M_z = \int_A E(ay + bz + c)y \, dA = Ea \int_A y^2 \, dA + Eb \int_A yz \, dA + Ec \int_A y \, dA$$
$$\Rightarrow M_z = Ea I_{zz} + Eb I_{yz} \quad - (2)$$

'a' and 'b' unknowns – solve simultaneous algebraic equations (1) and (2)

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Bending of arbitrary cross section beam

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- Solving (1) and (2) 'a' and 'b' values are as following

$$a = \frac{M_y I_{yz} - M_z I_{yy}}{E(I_{yz}^2 - I_{yy} I_{zz})}, \quad b = \frac{-M_y I_{zz} + M_z I_{yz}}{E(I_{yz}^2 - I_{yy} I_{zz})}$$

Strain is

$$\epsilon_{xx} = \frac{M_y I_{yz} - M_z I_{yy}}{E(I_{yz}^2 - I_{yy} I_{zz})} y + \frac{-M_y I_{zz} + M_z I_{yz}}{E(I_{yz}^2 - I_{yy} I_{zz})} z + \frac{N_x}{EA}$$

This represents a straight line

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Bending of arbitrary cross section beam

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- Bending stress – multiply axial strain (bending strain) with Young's modulus (E)
- Neutral axis /plane – no stress/strain => '0'

$$\frac{M_y I_{yz} - M_z I_{yy}}{(I_{yz}^2 - I_{yy} I_{zz})} y + \frac{-M_y I_{zz} + M_z I_{yz}}{(I_{yz}^2 - I_{yy} I_{zz})} z + \frac{N_x}{A} = 0$$

$$\Rightarrow z = \left(\frac{M_y I_{yz} - M_z I_{yy}}{M_y I_{zz} - M_z I_{yz}} \right) y - \frac{N_x}{A} (I_{yz}^2 - I_{yy} I_{zz})$$

$$\text{slope} \Rightarrow m = \left(\frac{M_y I_{yz} - M_z I_{yy}}{M_y I_{zz} - M_z I_{yz}} \right)$$

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Bending of arbitrary cross section beam

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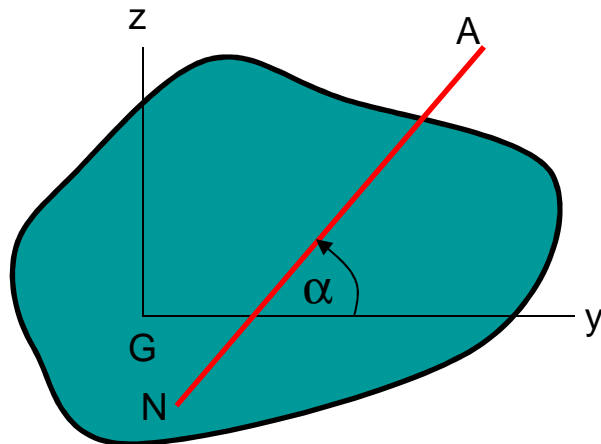
Equation of neutral axis

$$z = \left(\frac{M_y I_{yz} - M_z I_{yy}}{M_y I_{zz} - M_z I_{yz}} \right) y - \frac{N_x}{A} (I_{yz}^2 - I_{yy} I_{zz})$$

Neutral axis doesn't pass through origin – centroid of cross-section

If no axial load $\Rightarrow N_x = 0$

N.A equation $z = m y$ Passes through centroid



$$\text{slope, } m = \tan \alpha = \frac{M_y I_{yz} - M_z I_{yy}}{M_y I_{zz} - M_z I_{yz}}$$

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Bending of arbitrary cross section beam

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- Neutral axis (N. A)
 - Equation of neutral axis completely depends on geometry and loading – moments and axial load
 - Axial load offset N. A. – does not pass through centroid of cross-section
 - Orientation (slope) purely depends on moments and geometry
 - Orientation not decided by axial load
 - Independent of material properties - isotropic
 - Simplifying equation if co-ordinate system coincides with principal axes

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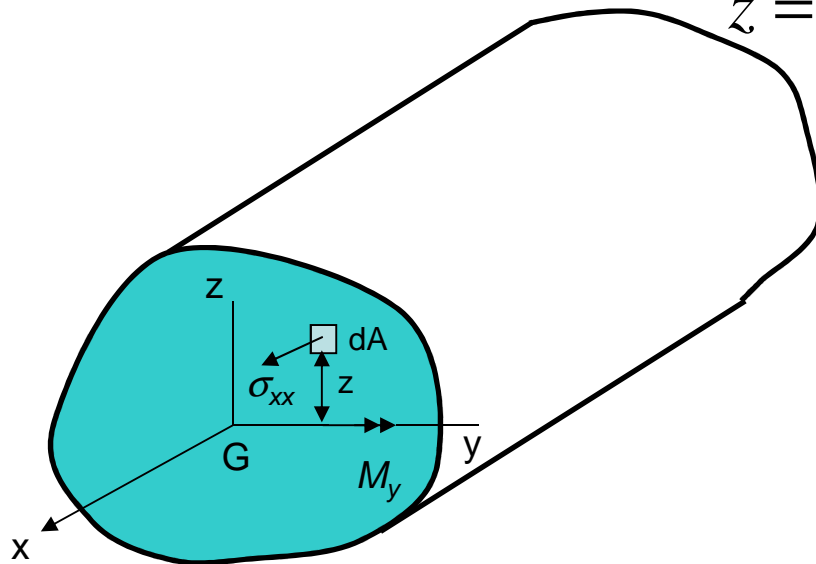


Bending of arbitrary cross section beam

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- Bending in a single plane x - z – no axial load

$$\Rightarrow N_x = 0$$



$$z = \left(\frac{M_y I_{yz} - M_z I_{yy}}{M_y I_{zz} - M_z I_{yz}} \right) y - \frac{N_x}{A} (I_{yz}^2 - I_{yy} I_{zz})$$

$$z = \left(\frac{M_y I_{yz} - M_z I_{yy}}{M_y I_{zz} - M_z I_{yz}} \right) y$$

$$\Rightarrow z = \frac{M_y I_{yz}}{M_y I_{zz}} y \Rightarrow z = \frac{I_{yz}}{I_{zz}} y$$

$$M_z = 0$$

Equation of N.A depends on area moment of inertia

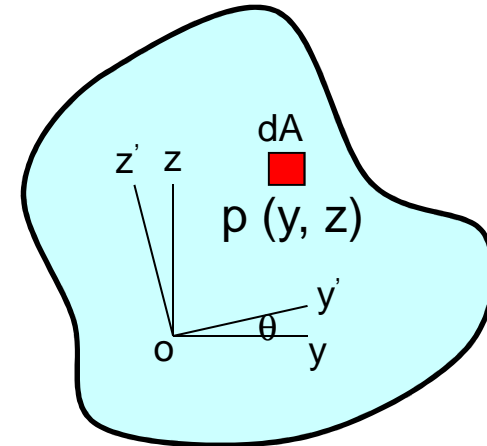


Principal axes

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- Area moment of inertia – geometric property

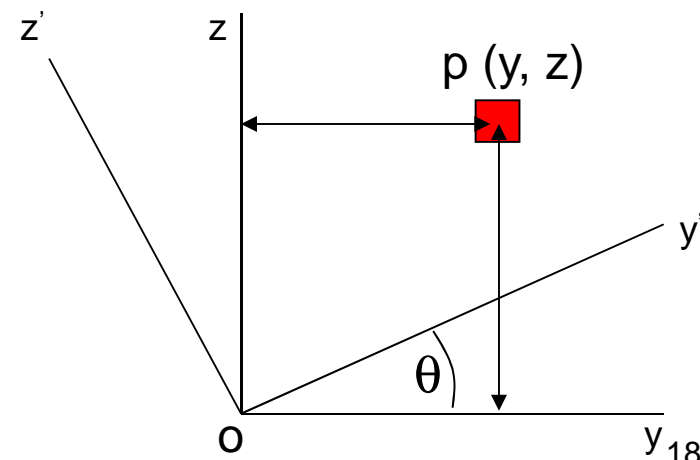
$$I_{yy} = \int_A z^2 dA, \quad I_{zz} = \int_A y^2 dA$$



Product moment of inertia

$$I_{yz} = \int_A yz dA$$

Rotate yOz csys to a new csys $y'Oz'$ at an angle ' θ '



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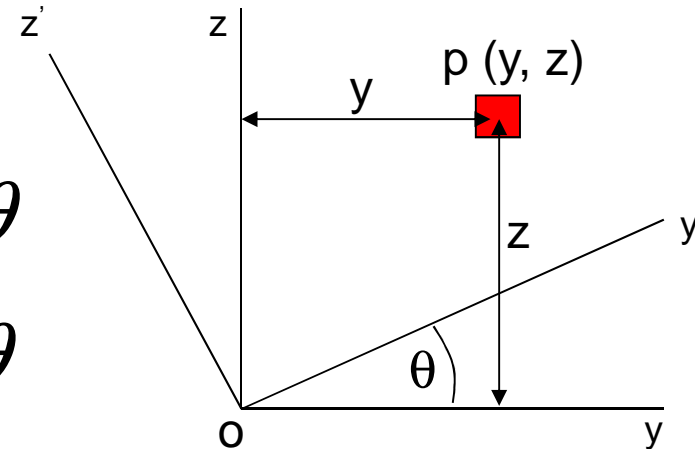
Principal axes

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Transformation

$$y = y' \cos \theta - z' \sin \theta$$

$$z = y' \sin \theta + z' \cos \theta$$



$$I_{zz} = \int_A y^2 dA = \int_A (y' \cos \theta - z' \sin \theta)^2 dA$$

$$I_{zz} = \cos^2 \theta \int_A y'^2 dA + \sin^2 \theta \int_A z'^2 dA - \sin 2\theta \int_A y' z' dA$$

$$I_{zz} = I_{z'z'} \cos^2 \theta + I_{y'y'} \sin^2 \theta - I_{y'z'} \sin 2\theta - (1)$$



Principal axes

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■ Transformation

$$I_{yy} = \int_A z^2 dA = \int_A (y' \sin \theta + z' \cos \theta)^2 dA$$

$$I_{yy} = \sin^2 \theta \int_A y'^2 dA + \cos^2 \theta \int_A z'^2 dA + \sin 2\theta \int_A y' z' dA$$

$$I_{yy} = I_{z'z'} \sin^2 \theta + I_{y'y'} \cos^2 \theta + I_{z'y'} \sin 2\theta - (2)$$



Principal axes

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- Product moment of inertia – a property defined wrt a set of perpendicular axes lying in the same plane of area

$$I_{yz} = \int_A yz dA$$

$$\Rightarrow I_{yz} = \int_A (y' \cos \theta - z' \sin \theta)(y' \sin \theta + z' \cos \theta) dA$$

$$\Rightarrow I_{yz} = \cos \theta \sin \theta \left(\int_A y'^2 dA - \int_A z'^2 dA \right) + (\cos^2 \theta - \sin^2 \theta) \int_A y' z' dA$$



Principal axes

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■ Product M. I

$$I_{yz} = \frac{1}{2} \sin 2\theta (I_{z'z'} - I_{y'y'}) + \cos 2\theta I_{y'z'} \quad - (3)$$

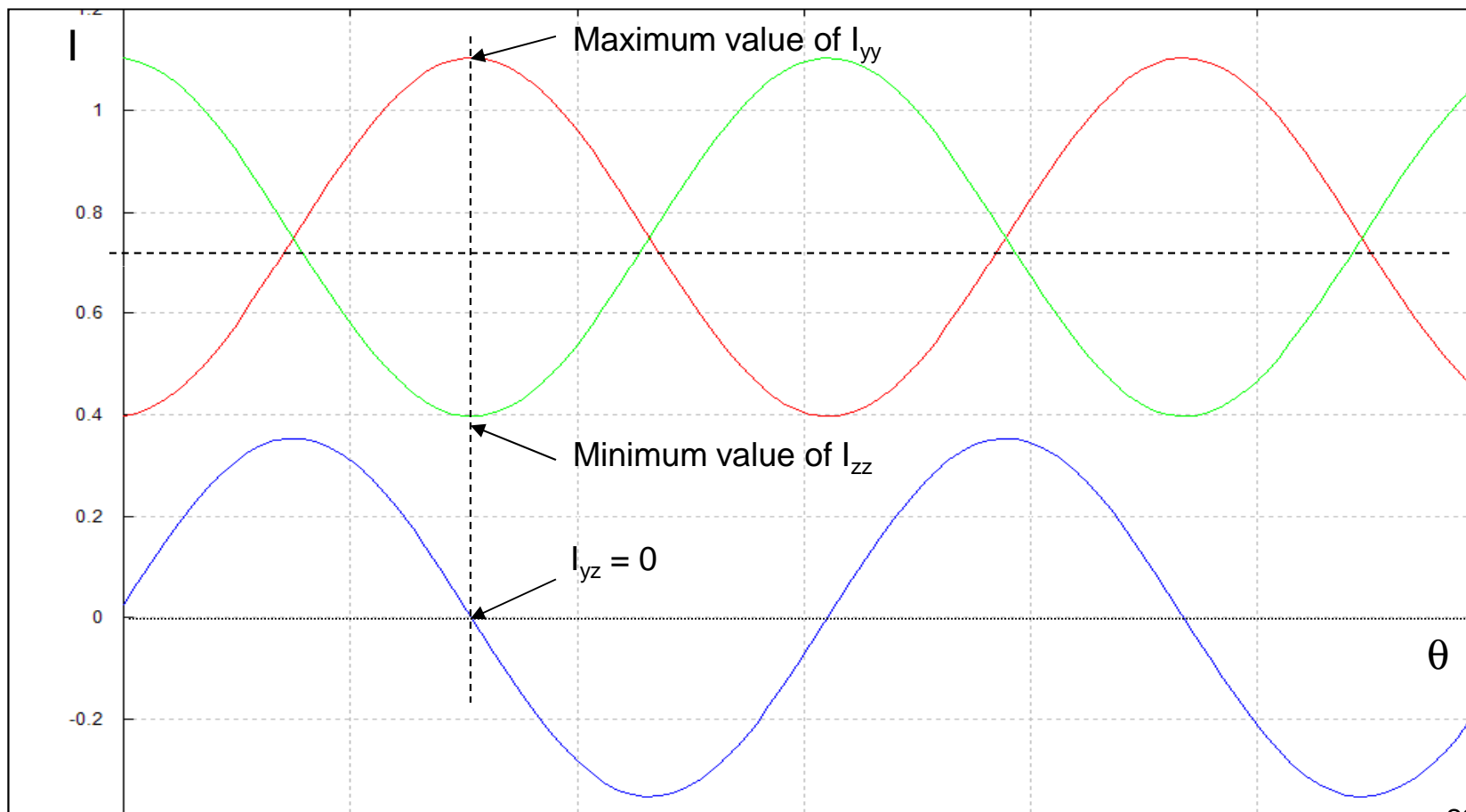
- Product M. I can be +ve or -ve – depends on selection of co-ordinate system
- Variation of I_{yy} , I_{zz} and I_{yz} wrt 'θ' – harmonic
- When product M. I changes sign – it passes through zero also
- Principal axes is a csys, where product M.I is zero



Principal axes

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- Variation of I_{yy} , I_{zz} and I_{yz}



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Principal axes

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- Product M. I = 0, area M. I – maximum and minimum
- Orientation corresponding to max. & min. – principal axes $\Rightarrow \theta, \theta + 90^\circ$
- From (3), $I_{yz} = 0$

$$I_{yz} = \frac{1}{2} \sin 2\theta (I_{z'z'} - I_{y'y'}) + \cos 2\theta I_{y'z'} = 0$$

$$\Rightarrow \tan 2\theta = \frac{2I_{y'z'}}{(I_{y'y'} - I_{z'z'})}$$

Area moment of inertia – a second order tensor

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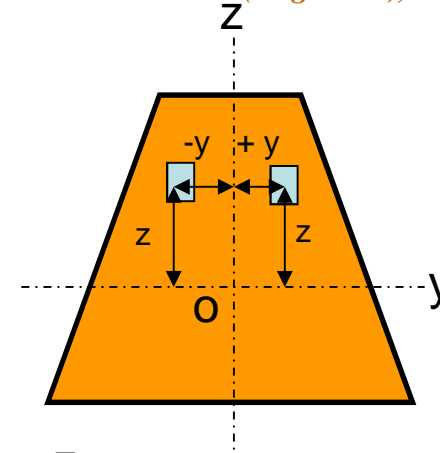
Principal axes

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■ Symmetric sections

Product moment of inertia

$$I_{yz} = \int_A yz dA$$

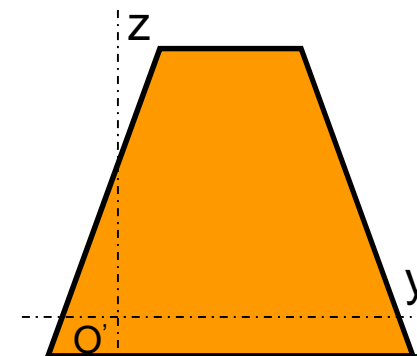


Co-ordinate system is selected symmetrically, I_{yz} is positive and negative.

Summation over the area vanishes zero. YoZ – principal co-ordinate system.

$Yo'Z$ – another co-ordinate system.
Product moment of inertia doesn't vanish.

One axis of symmetry – principal axis



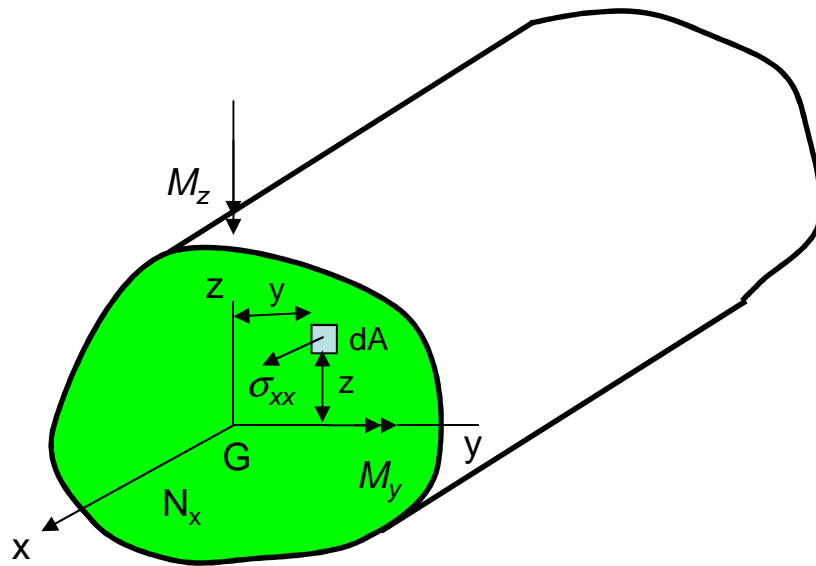
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Bending of arbitrary cross section beam

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- Co-ordinate system is aligned with principal csys – product M. I $\Rightarrow I_{yz} = 0$



$$\epsilon_{xx} = \frac{M_y I_{yz} - M_z I_{yy}}{E(I_{yz}^2 - I_{yy} I_{zz})} y + \frac{-M_y I_{zz} + M_z I_{yz}}{E(I_{yz}^2 - I_{yy} I_{zz})} z + \frac{N_x}{EA}$$

$$\epsilon_{xx} = \frac{M_z}{EI_{zz}} y + \frac{M_y}{EI_{yy}} z + \frac{N_x}{EA}$$

Simplified expression

Again, N. A equation can be obtained by making strain zero

$$y = - \frac{M_y I_{zz}}{M_z I_{yy}} z - \frac{N_x I_{zz}}{A M_z}$$

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Bending of arbitrary cross section beam

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- When no axial load acts – $N_x = 0$ –
equation of N. A

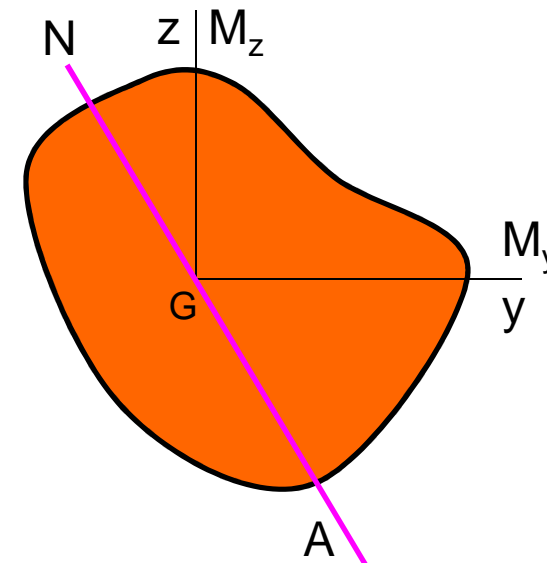
$$y = -\frac{M_y}{M_z} \frac{I_{zz}}{I_{yy}} z - \frac{N_x}{A} \frac{I_{zz}}{M_z}$$

$$\Rightarrow y = -\frac{M_y}{M_z} \frac{I_{zz}}{I_{yy}} z$$

Straight line passing through centroid

If, $M_z = 0 \Rightarrow$ Equation of N. A $\Rightarrow z = 0$

This is 'y' axis

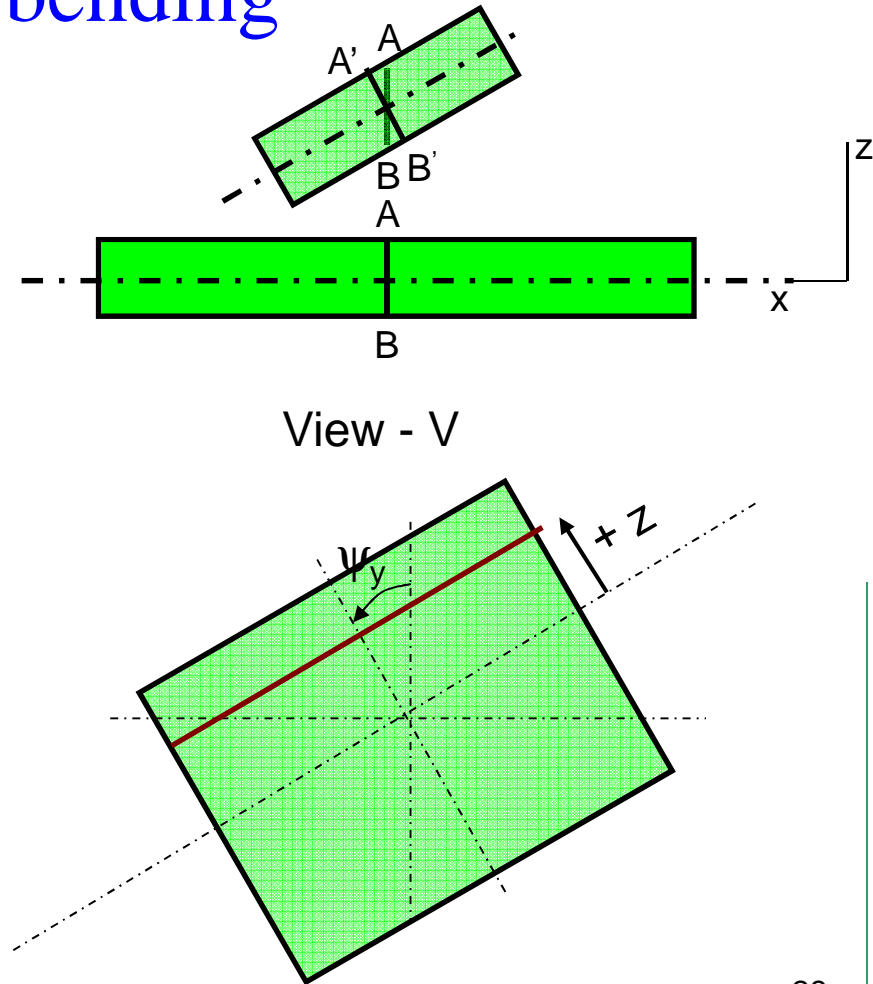
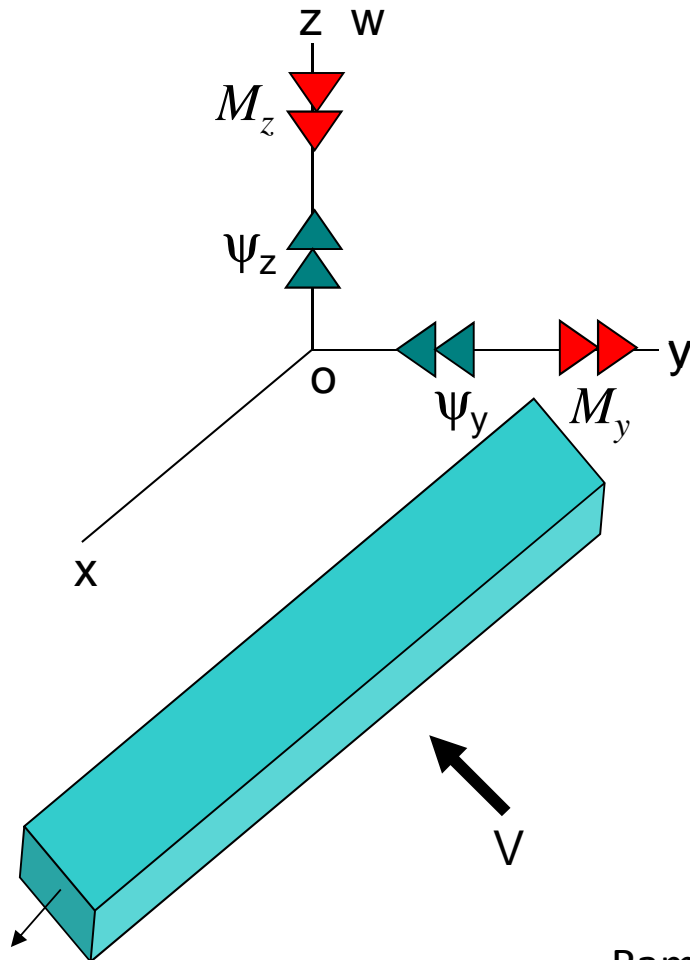




Axial deformation

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■ Axial deformation in bending



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Axial deformation

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- Axial deformation due to axial load,
 $N_x \Rightarrow u_o$
- Axial deformation due to bending
 - Bending in xz plane $\Rightarrow -z\psi_y$
 - Bending in xy plane $\Rightarrow -y\psi_z$
- Axial deformation due to all loads

$$u(x) = u_o - z\psi_y(x) - y\psi_z(x)$$

-ve signs – rotations are opposite to moments

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Axial deformation

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- In Euler-Bernoulli's theory – shears strains negligible

$$u = u_o - z\psi_y - y\psi_z$$

$$\gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} = 0$$

$$\Rightarrow -\psi_z + \frac{\partial v}{\partial x} = 0$$

$$\Rightarrow \frac{\partial v}{\partial x} = \psi_z$$

$$\gamma_{xz} = \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} = 0$$

$$\Rightarrow -\psi_y + \frac{\partial w}{\partial x} = 0$$

$$\Rightarrow \frac{\partial w}{\partial x} = \psi_y$$



Axial deformation

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- Axial deformation in terms of deflection gradient

$$u = u_o - z \frac{\partial w}{\partial x} - y \frac{\partial v}{\partial x}$$

$$\text{strain, } \epsilon_{xx} = \frac{\partial u}{\partial x} = \frac{\partial u_o}{\partial x} - z \frac{\partial^2 w}{\partial x^2} - y \frac{\partial^2 v}{\partial x^2}$$

$$\epsilon_{xx} = c + bz + ay$$

$$c = \frac{\partial u_o}{\partial x}, \quad b = -\frac{\partial^2 w}{\partial x^2}, \quad a = -\frac{\partial^2 v}{\partial x^2}$$



Curvatures

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- Curvature – gradient of slope

$$a = -\frac{\partial^2 v}{\partial x^2} = \frac{M_z}{EI_{zz}} \Rightarrow -E \frac{\partial^2 v}{\partial x^2} = \frac{M_z}{I_{zz}}$$

$$b = -\frac{\partial^2 w}{\partial x^2} = \frac{M_y}{EI_{yy}} \Rightarrow -E \frac{\partial^2 w}{\partial x^2} = \frac{M_y}{I_{yy}}$$

$$c = \frac{\partial u_o}{\partial x} = \frac{N_x}{EA} \Rightarrow EA \frac{\partial u_o}{\partial x} = N_x$$

Moment – curvature relationships

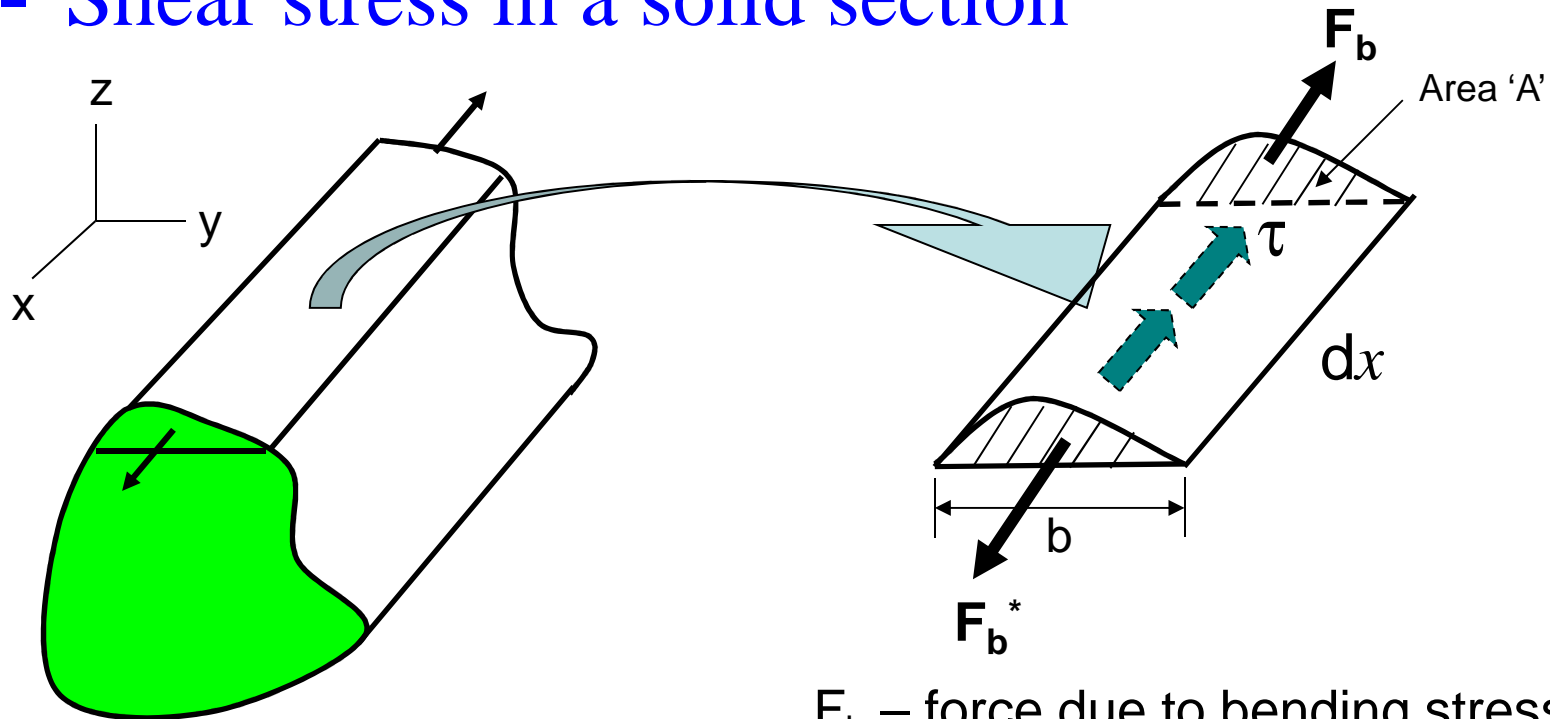
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Shear stress distribution

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■ Shear stress in a solid section



Solid section

F_b – force due to bending stress

τ - shear acting on plane $b \, dx$

'b' – local width on which shear to be estimated

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Shear stress distribution

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- Force due to bending stress on area 'A'

$$F_b = \int_A \sigma_{xx} dA$$

Bending force at $x+dx$

$$F_b^* = F_b + \frac{\partial F_b}{\partial x} dx = \int_A \sigma_{xx} dA + \int_A \frac{\partial \sigma_{xx}}{\partial x} dA$$

Shear force acting on plane $b dx$

$$F_s = \tau b dx$$

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Shear stress distribution

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Equilibrium

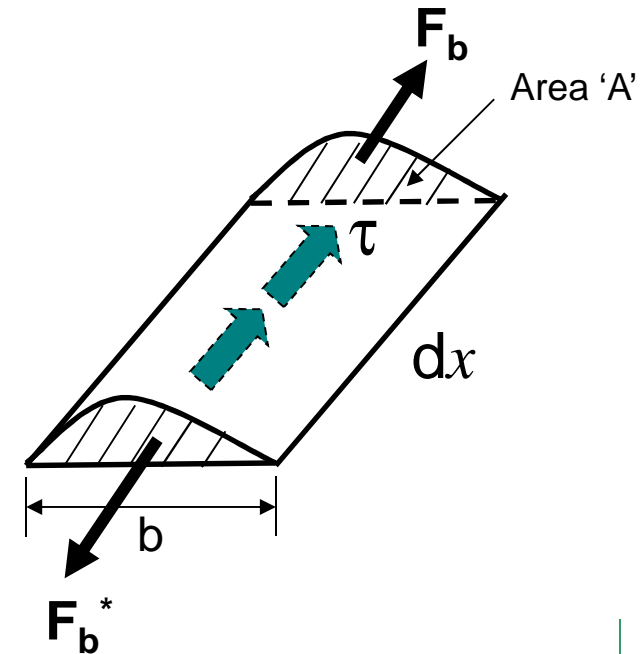
$$F_b^* - F_b - F_s = 0$$

$$F_b + \frac{\partial F_b}{\partial x} dx - F_b - \tau b dx = 0$$

$$\Rightarrow \frac{\partial F_b}{\partial x} = \tau b$$

$$F_b = \int_A \sigma_{xx} dA$$

$$\Rightarrow \frac{\partial}{\partial x} \int_A \sigma_{xx} dA = \tau b \Rightarrow \tau b = \int_A \frac{\partial \sigma_{xx}}{\partial x} dA$$



A = area, which is above the plane on which shear to be found

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Shear stress distribution

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■ Bending stress

$$\sigma_{xx} = E \varepsilon_{xx} = \frac{M_z}{I_{zz}} y + \frac{M_y}{I_{yy}} z + \frac{N_x}{A}$$

$$\frac{\partial \sigma_{xx}}{\partial x} = \frac{1}{I_{zz}} \frac{\partial M_z}{\partial x} y + \frac{1}{I_{yy}} \frac{\partial M_y}{\partial x} z$$

If a beam is subjected to pure bending – no variation of bending moment along length

$$\frac{\partial M_z}{\partial x} = 0, \quad \frac{\partial M_y}{\partial x} = 0, \quad \frac{\partial \sigma_{xx}}{\partial x} = 0 \Rightarrow \tau = 0$$

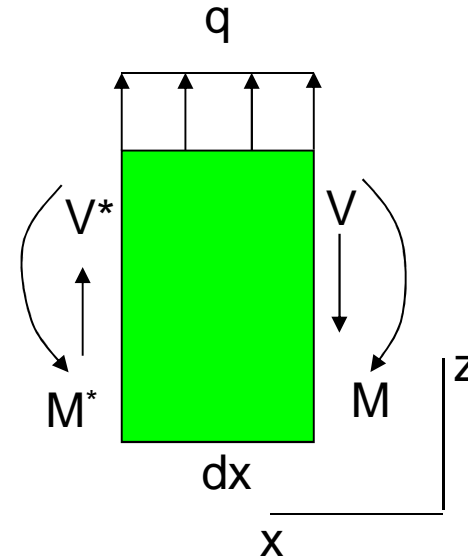
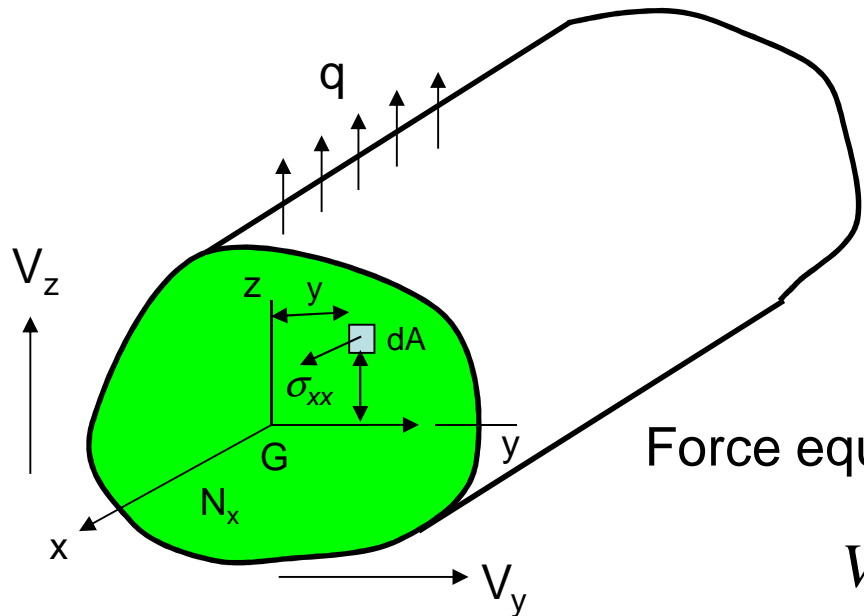
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Shear stress distribution

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- Shear stress exists when transverse loads are applied



Force equilibrium of infinitesimal element

$$V^* + qdx - V = 0$$

$$\Rightarrow V + \frac{\partial V}{\partial x} dx + qdx - V = 0 \Rightarrow \frac{\partial V}{\partial x} = -q \quad (1)$$

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Shear stress distribution

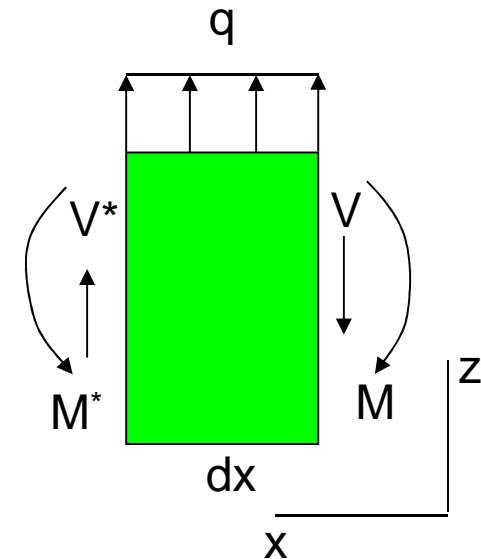
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■ Moment equilibrium

$$M^* - M - V^* dx - q dx \frac{dx}{2} = 0$$

$$M + \frac{\partial M}{\partial x} dx - M - \left(V + \frac{\partial V}{\partial x} dx \right) dx - q \frac{(dx)^2}{2} = 0$$

$$\frac{\partial M}{\partial x} = V \quad - (2)$$



Higher order terms are neglected

From (1) and (2)

$$\frac{\partial}{\partial x} \left(\frac{\partial M}{\partial x} \right) = \frac{\partial V}{\partial x} = -q \Rightarrow \frac{\partial^2 M}{\partial x^2} = -q$$

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Shear stress distribution

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- When transverse loads are acting

Using equation (2) $V_z = \frac{\partial M_y}{\partial x}$, $V_y = \frac{\partial M_z}{\partial x}$

$$\tau b = \int_A \frac{\partial \sigma_{xx}}{\partial x} dA$$

$$\frac{\partial \sigma_{xx}}{\partial x} = \frac{1}{I_{zz}} \frac{\partial M_z}{\partial x} y + \frac{1}{I_{yy}} \frac{\partial M_y}{\partial x} z$$

$$\frac{\partial \sigma_{xx}}{\partial x} = \frac{V_y}{I_{zz}} y + \frac{V_z}{I_{yy}} z$$

$$\tau b = \int_A \left(\frac{V_y}{I_{zz}} y + \frac{V_z}{I_{yy}} z \right) dA$$

$$A \bar{y} = \int_A y dA, \quad A \bar{z} = \int_A z dA$$

$$\tau b = \frac{V_y}{I_{zz}} A \bar{y} + \frac{V_z}{I_{yy}} A \bar{z}$$



Shear stress distribution

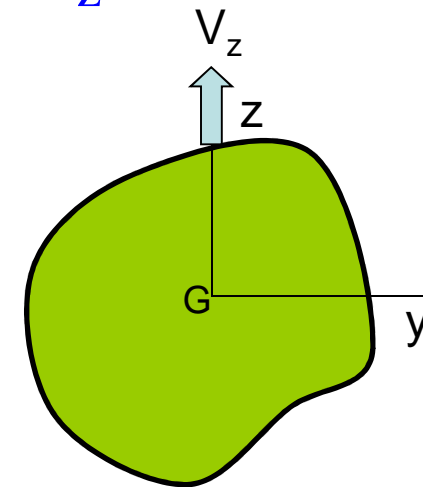
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- Transverse loading in xz plane - V_z

$$\tau b = \frac{V_y}{I_{zz}} A \bar{y} + \frac{V_z}{I_{yy}} A \bar{z} \quad V_y = 0$$

$$\tau b = \frac{V_z}{I_{yy}} A \bar{z}$$

Resembles shear flow, q



Shear flow at a given horizontal section depends on moment of area above that section wrt csys passing through centroid

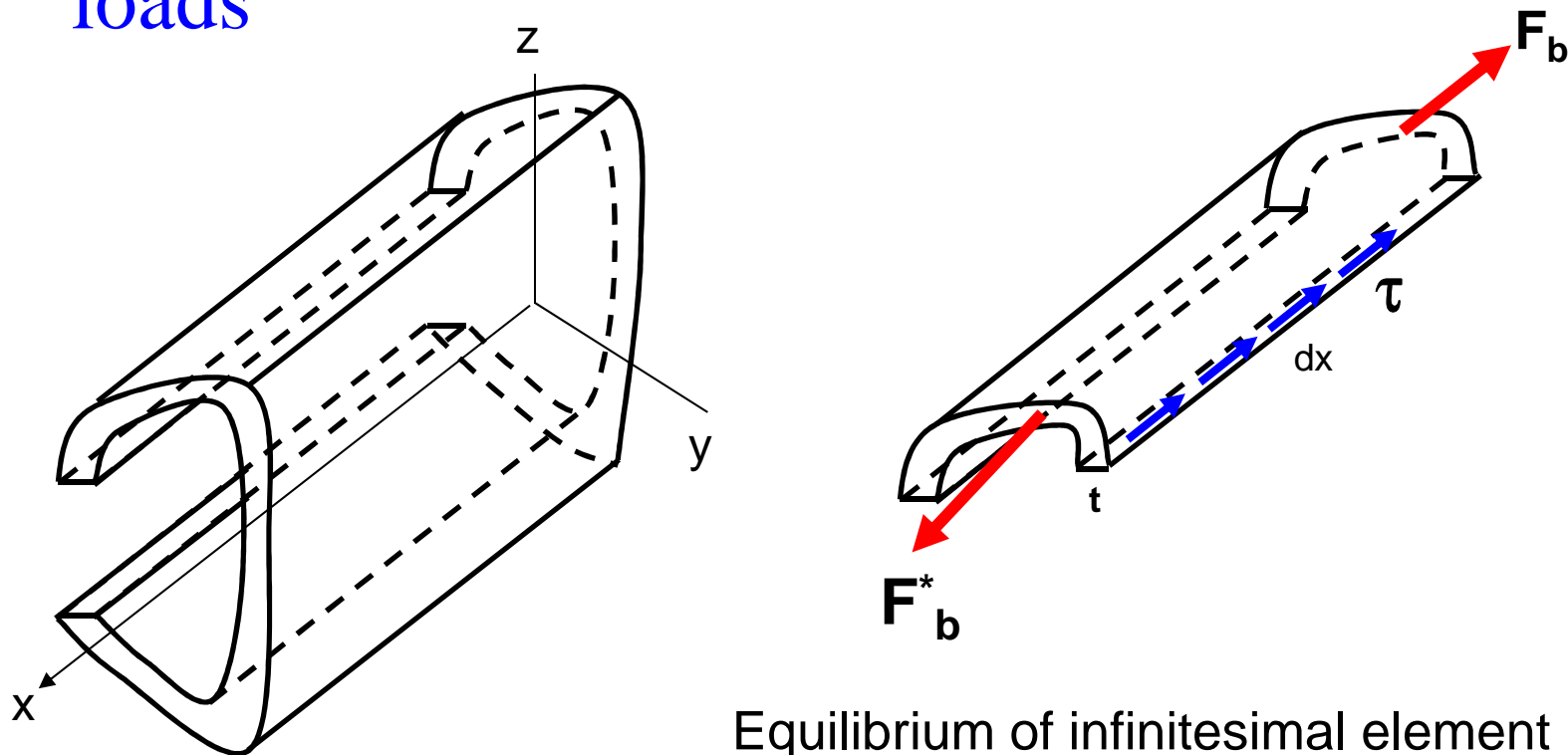
This shear flow is not constant. Varies from zero to maximum



Shear stress distribution

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- Thin walled beam subjected to transverse loads



Equilibrium of infinitesimal element gives stress distribution in the wall

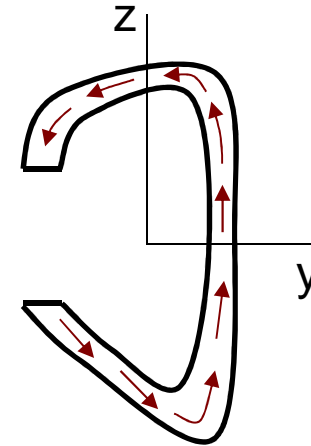


Shear stress distribution

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■ Shear stress in thin walled beam

$$\tau t = \frac{V_y}{I_{zz}} A \bar{y} + \frac{V_z}{I_{yy}} A \bar{z} = q$$



Shear flow in the wall of thin beam

Sum of the shear forces in the cross-section = externally applied shear force => vertical equilibrium of the section

Shear flow is not constant – varies from point to point along the wall

Shear is purely because of transverse loads, which cause bending => Bending – shear

Shear due to pure torsion => Torsion - shear
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Shear center

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- Equilibrium of beam cross-section requires
 - Forces and moment equilibrium

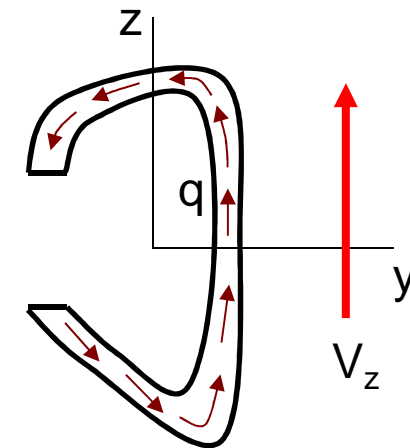
Twisting will not take place if moment due to shear developed in the wall is equal to moment due to external transverse load

Variable – location of application of external load

Shear flow is a function of external transversal load.

Moments of forces taken wrt any point in space – gives point of application of external load – shear center

Shear center – independent of external load – depends only on geometry

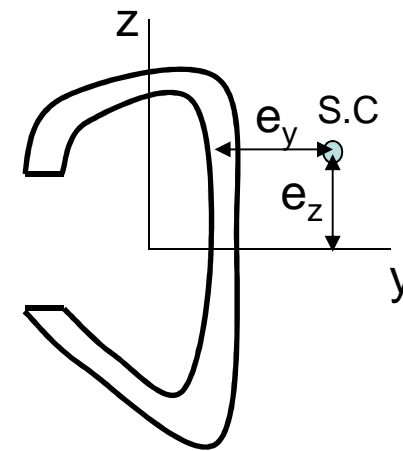
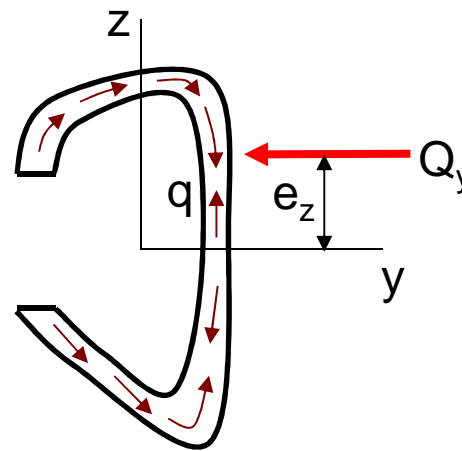
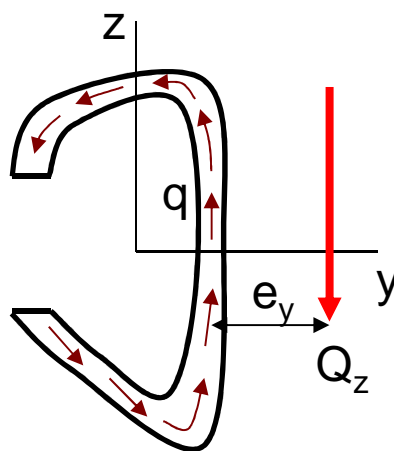




Shear center

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- Shear center (S. C) may or may not pass through centroid of beam cross-section
- When there is no axial load, N. A passes through centroid
- Loads acting in both planes



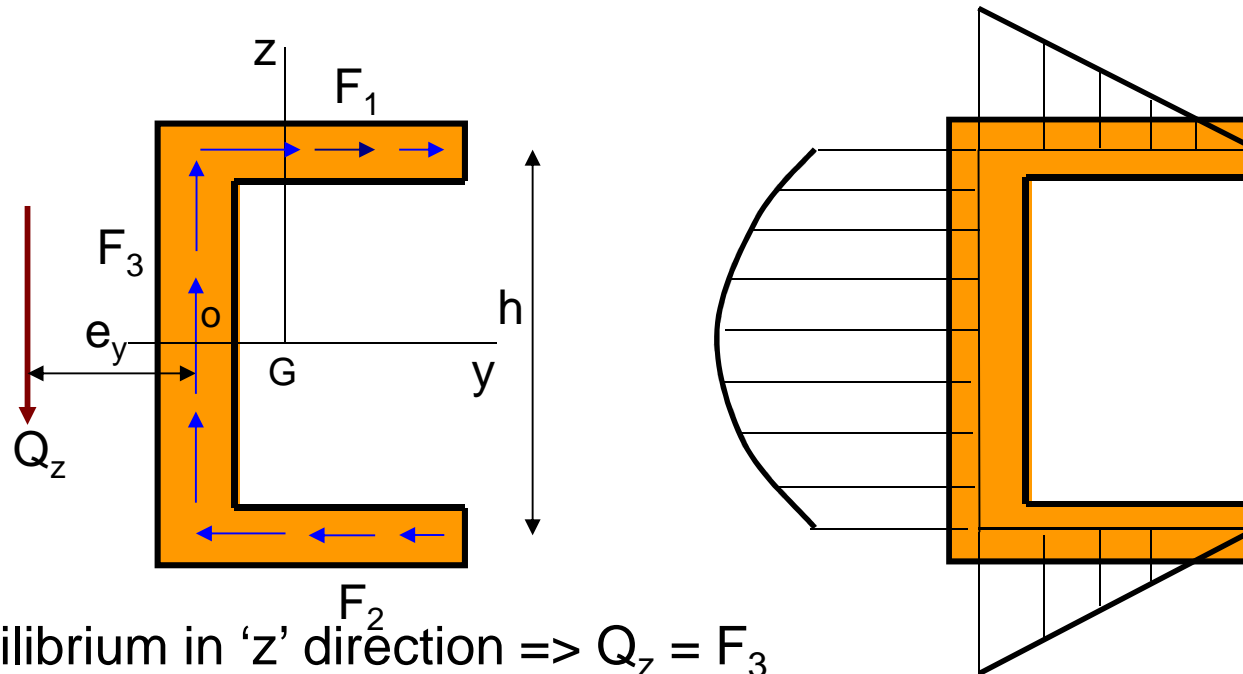
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Shear center

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■ Some specific cross-sections



Equilibrium in 'z' direction $\Rightarrow Q_z = F_3$

Equilibrium in 'y' direction $\Rightarrow F_1 - F_2 = 0$

Moment of forces wrt 'o' $Q_z e_y - F_1 \frac{h}{2} - F_2 \frac{h}{2} = 0 \Rightarrow e_y = \frac{F_1}{Q_z} h$

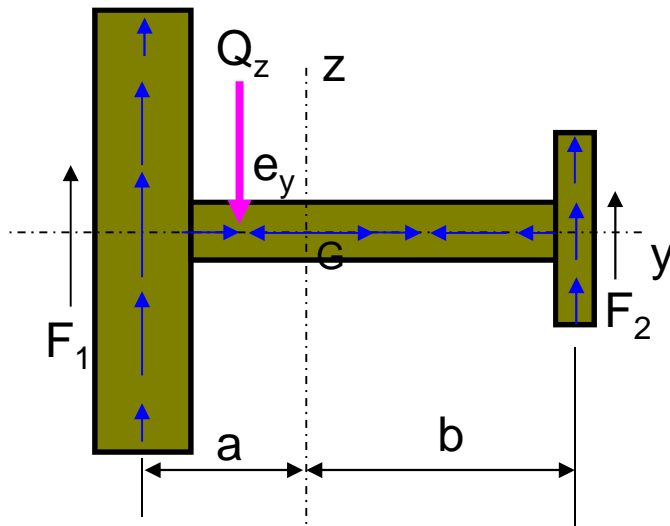
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Shear center

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- When load acts at S.C, it produces only bending – no twisting takes place



Vertical equilibrium $\Rightarrow F_1 + F_2 = Q_z$

Horizontal equilibrium – sum of forces in horizontal web = 0

Taking moments wrt 'G'

$$F_1 a - F_2 b - Q_z e_y = 0$$

$$\Rightarrow e_y = \frac{F_1 a - F_2 b}{Q_z}$$

Express F_1 and F_2 in terms of transverse load, Q_z

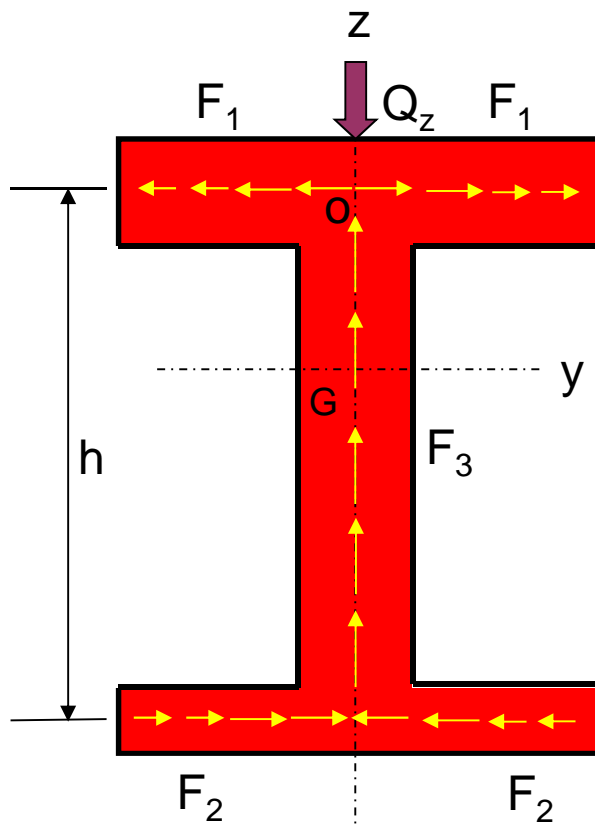
e_y will be a function of geometric parameters



Shear center

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■ Symmetric cross-section



Since cross-section is symmetric, shear distribution in each leg is symmetric wrt xz plane

Taking moments about 'O'

$$F_2 h - F_2 h = 0 \quad - \text{no moment}$$

External load should pass through axis of symmetry – shear center lies in the plane symmetry

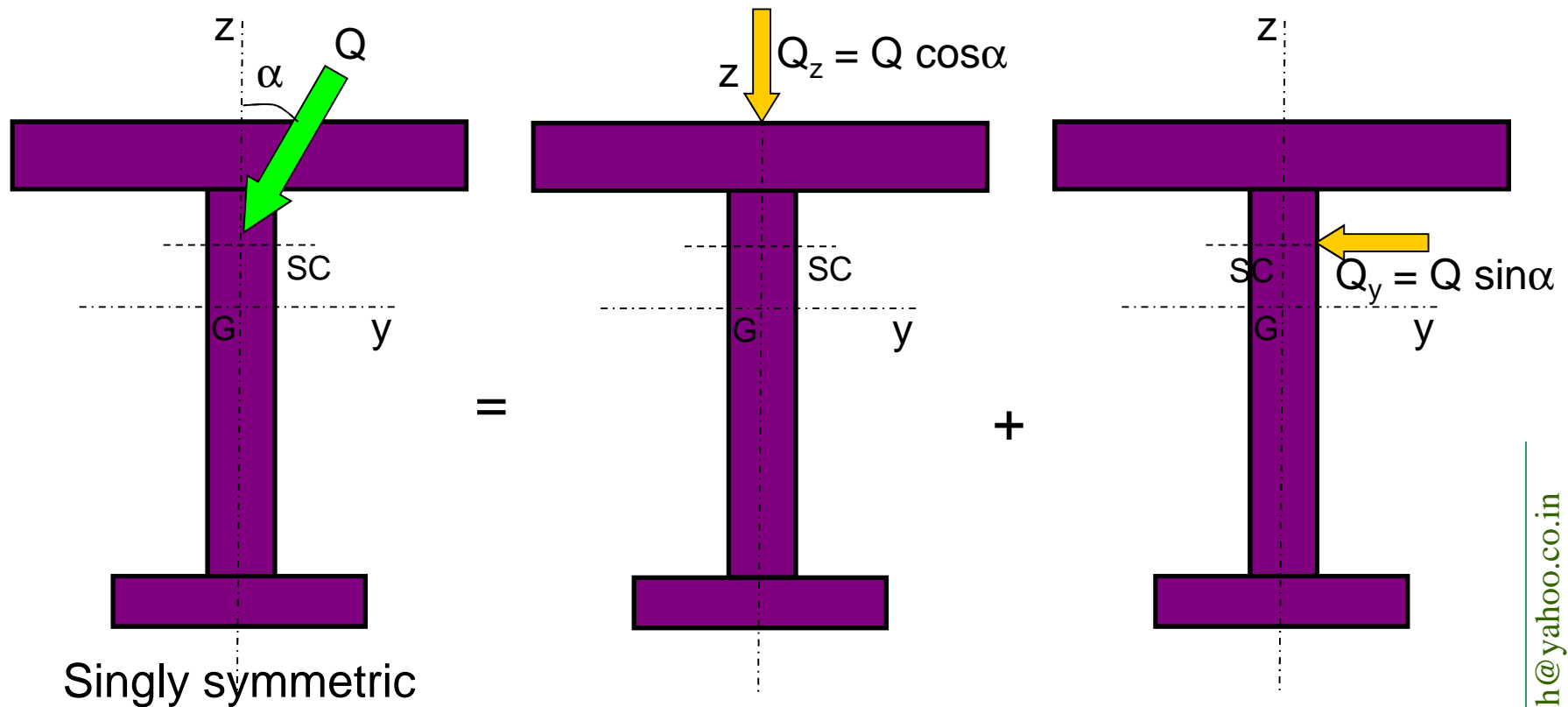
In symmetric cross-sections shear center lies in the plane of symmetry



Shear center

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- Load acting at some angle



Singly symmetric

Point of application of load has to through S.C
– no twisting takes place

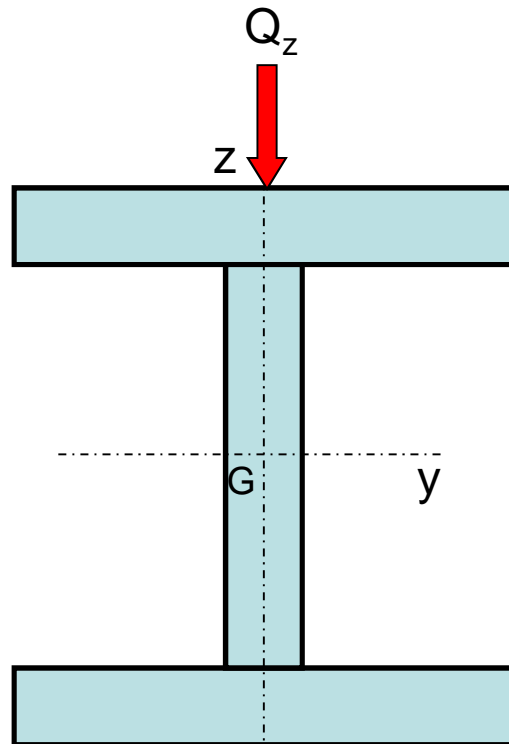
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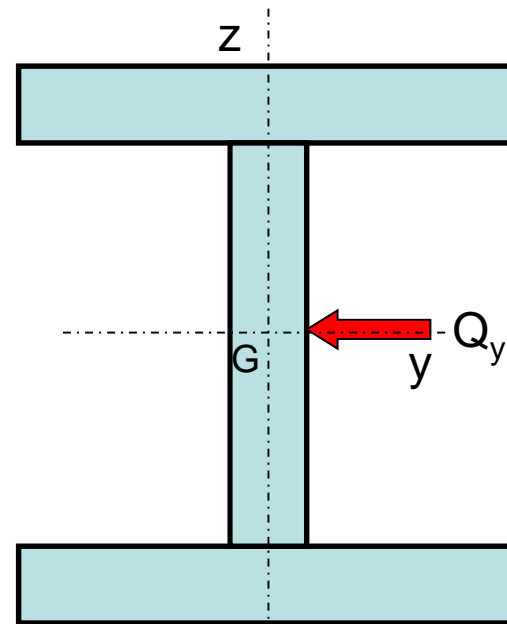
Shear center

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■ Doubly symmetric beam



Symmetric in xz plane – S.C
lies on z - axis



Symmetric in xy plane – S.C
lies on y - axis

In a doubly
symmetric
beam shear
center lies
at centroid

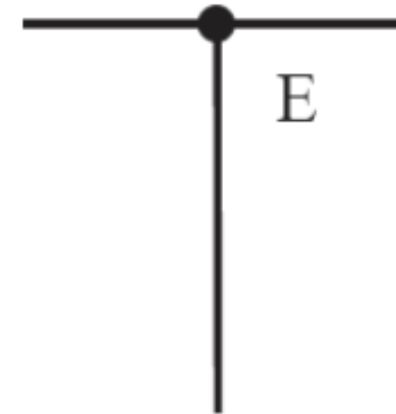
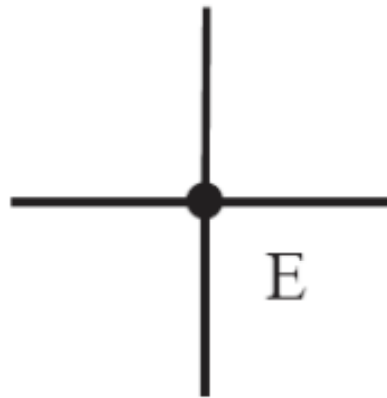
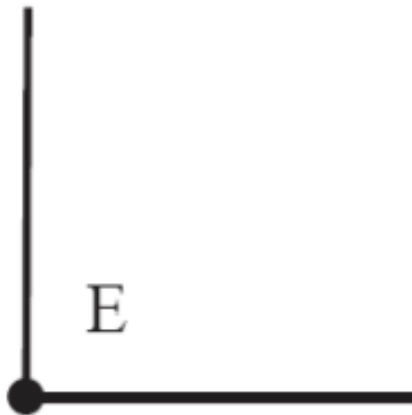
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Location of shear center for some cross-sections





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